Nuclear Anapole Moments and Their Constraints on the PNC Nuclear Interaction

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- What is the anapole moment?
- The PNC nuclear interaction and nuclear AMs
- Experimental and theoretical works
- Summary

- What is the anapole moment?
- **■** From classical EM:

$$\Phi(\vec{x}) = \int d^3 x' \frac{\rho(\vec{x}')}{4\pi |\vec{x} - \vec{x}'|}$$

$$= \int d^3 x' \rho(\vec{x}') \left\{ 1 - \vec{x}' \cdot \vec{\nabla} + \frac{1}{2} \left(\vec{x}' \cdot \vec{\nabla} \right)^2 + \cdots \right\} \frac{1}{4\pi |\vec{x}|}$$

$$\rightarrow \text{ net charge, charge dipole, charge quadrupole } \oplus \text{ charge radius,}$$

$$\overrightarrow{A}(\overrightarrow{x}) = \int d^3 x' \frac{\overrightarrow{j}(\overrightarrow{x}')}{4\pi |\overrightarrow{x} - \overrightarrow{x}'|}$$

$$= \int d^3 x' \overrightarrow{j}(\overrightarrow{x}') \left\{ 1 - \overrightarrow{x}' \cdot \overrightarrow{\nabla} + \frac{1}{2} \left(\overrightarrow{x}' \cdot \overrightarrow{\nabla} \right)^2 + \cdots \right\} \frac{1}{4\pi |\overrightarrow{x}|}$$

$$\rightarrow \text{ no net current, magnetic dipole, magnetic quadrupole } \oplus$$

anapole, ...

■ Anapole Moment:

$$\vec{a} = -\frac{1}{4} \int d^3 x' \ \vec{x}'^2 \vec{j}(\vec{x}')$$

$$\vec{A}(\vec{x}) = \vec{a} \, \delta^3(\vec{x})$$

- $\rightarrow \vec{a}$ is odd (even) for a vector (axial-vector) current
- $\rightarrow \vec{A}(\vec{x})$ is of contact form

■ Multipole Classification:

	C_J	E_J	M_J
J = 0	PT		
J=1			PT
J=2	PT	P	
J=3			PT
	•		

- $\rightarrow C_0$: charge monopole
- $\rightarrow C_1$: charge dipole
- $\rightarrow M_1$: magnetic dipole
- $\rightarrow E_1$: anapole

■ Definition:

- → The EM vector multipole which is P-odd but T-even
- → By multipole expansion:

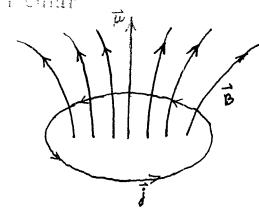
$$a_{\lambda} = \frac{-i\sqrt{6\pi}}{\vec{q}^2} T_{1\lambda}^{\text{el}}$$

 \rightarrow A dimensionless constant κ is often used:

$$\vec{a} = \frac{G_F}{\sqrt{2}} \kappa \vec{I} \text{ (nuclear spin)}$$

Look at the current windings

a Cicalar



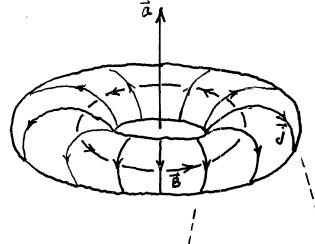
 $\vec{j}(\vec{x})$: toroidal

 $\vec{B}(\vec{x})$: poloidal (extensive)

 $\vec{\mu} \neq 0$

 $\vec{a} = 0$

■ Toroidal

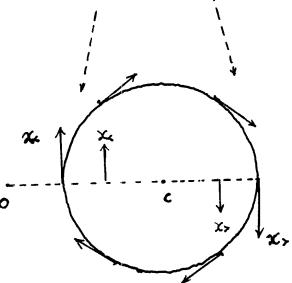


 $\vec{j}(\vec{x})$: poloidal

 $\vec{B}(\vec{x})$: toroidal (confined)

 $\overrightarrow{\mu} = 0$

 $\vec{a} \neq 0$



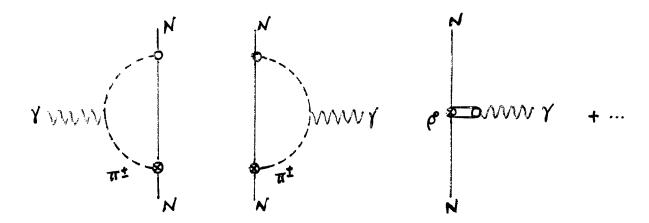
• The PNC nuclear interaction and nuclear AMs

PNC nuclear interaction is important to understand both the hadronic neutral weak interaction and dynamics of strong interaction

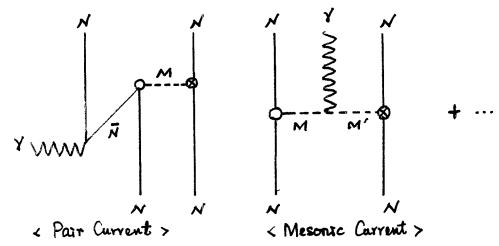
- Long-range PNC N-N interaction from the light meson exchange scheme
- → Direct W^{\pm} and Z^{0} exchanges is suppressed by the hard core (\lesssim 0.5 fm)
- → The physics of how weak bosons exchange between quarks at short distances is treated phenomenologically by a set of meson-nucleon PV couplings
- \rightarrow Conventional framework is OBEP with π^{\pm} , ρ , and ω as mediators, the six PNC meson-nucleon coupling constants are constrained by experiments while DDH best values provide the theoretical benchmark

Couplings	Equivalent	Best ($\times 10^{-7}$)	Range ($\times 10^{-7}$)
f_{π}	$\sqrt{32} F_{\pi}/g_{\pi NN}$	4.56	0 ~ 11.4
$h_{ ho}^{0}$	$-2F_0/g_{ ho NN}$	-11.4	-30.78 ~ 11.4
$h_{ ho}^{1}$	$-2F_1/g_{ ho NN}$	-0.19	-0.38 ~ 0
$h_{ ho}^2$	$-2F_2/g_{ ho { m NN}}$	-9.5	$-11.02 \sim -7.6$
h_{ω}^{0}	$-2G_0/g_{\omega \rm NN}$	-1.9	-10.26 ~ 5.7
h^1_ω	$-2G_1/g_{\omega NN}$	-1.14	-1.9 ~ -0.76

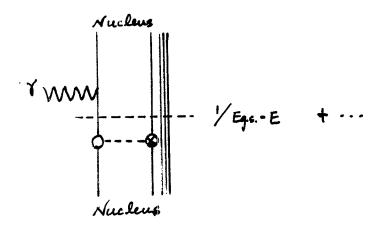
- PNC electromagnetic couplings induced by $H_{\text{PNC}}^{(2)}$
- → 1-body contribution (nucleonic AM)



→ 2-body contribution (PNC exchange currents)



→ Polarization mixing (E1 excitations)



Experimental and theoretical works

- Atomic PNC experiments
- → Asymmetry in highly forbidden M1 transitions

e.g.
$$6S_{1/2} \longleftrightarrow 7S_{1/2}$$
 in ¹³³Cs

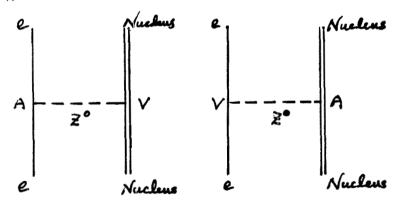
→ Optical rotation in allowed M1 transitions

e.g.
$$6P_{1/2} \longleftrightarrow 6P_{3/2}$$
 in 203,205 Tl

■ Contributing processes

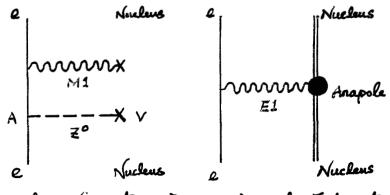
(1) Z^0 exchange

 \rightarrow NSI part ($\propto Q_W$) dominates, NSD part suppressed by $1 \rightarrow 1 \sin^2 \theta_W$



(2) NSD radiative corrections of $O(G_F \alpha)$

 \rightarrow AM gives the dominant NSD effect for heavy nuclei ($\sim A^{2/3}$)



1 Huperline Correction > 1 Anapale Interaction

Extraction of anapole contribution

$$H_{\mathrm{PNC}}^{\mathrm{NSD}}(e - \mathrm{nucleus}) = \frac{G_F}{2} \kappa_{\mathrm{tot}} \, \vec{\alpha} \cdot \vec{l} \, \rho(\vec{r})$$

 $\kappa_{\mathrm{tot}} = \kappa_Z + \kappa_{\mathrm{hf}} + \kappa_{\mathrm{anapole}}$

→ Careful experiments, and good atomic many-body calculations, get:

$$\kappa_{\text{tot}}(^{133}\text{Cs}) = 0.112 \pm 0.016 \text{ (1997, Colorado)}$$
 $\kappa_{\text{tot}}(^{205}\text{Tl}) = 0.293 \pm 0.400 \text{ (1995, Seattle)}$
 $\kappa_{\text{tot}}(^{205}\text{Tl}) = -0.08 \pm 0.40 \text{ (1995, Oxford)}$

→ Subtraction of Z exchange and hyperfine correction, get:

$$\kappa_{\rm anapole}(^{133}{\rm Cs}) = 0.090 \pm 0.016$$
 (7 σ determination)
 $\kappa_{\rm anapole}(^{205}{\rm Tl}) = 0.376 \pm 0.400$ (consistent with 0)

- Ongoing works and proposals on atomic PNC
- → Cs (Paris)
- → Tl (Seattle): sign and value?
- → Dy, Yb (Berkeley): odd neutron
- → Fr (Stony Brook, TRIUMF): isotope effect
- \rightarrow NSI effect measures Q_W , NSD effect shows nuclear AM!

■ Theoretical calculations of nuclear AMs

$$\langle \Psi \mid T_1^{\text{el}} \mid \Psi \rangle = \langle \Psi_0 \mid T_1^{\text{el}(A)} \mid \Psi_0 \rangle$$

$$+\sum_{\Phi_0} \left(\frac{\langle \Psi_0 | T_1^{\mathrm{el}(V)} | \Phi_0 \rangle \langle \Phi_0 | H_{\mathrm{PNC}}^{(2)} | \Psi_0 \rangle}{E_{\Psi_0} - E_{\Phi_0}} + H.c. \right)$$

where all states are many-body states

- → Subtleties:
- 1. The form of H_{PNC} : 2-body or 1-bodyanalog
- 2. Current conservation: should include PNC meson exchange currents arising from $H_{\rm PNC}^{(2)}$, and note that some of them might not be constrained by the gauge principle
- 3. Good description of nuclear structure: one of the fundamental problem in nuclear many-body physics
- Selected results of κ_{anapole} using DDH best values

A: s.p. approximation

B: s.p.

Woods-Saxon, oscillator, full currents ...etc.

C: RPA with constant density approximation

D: RPA beyond const. density

E: SM with closure approximation

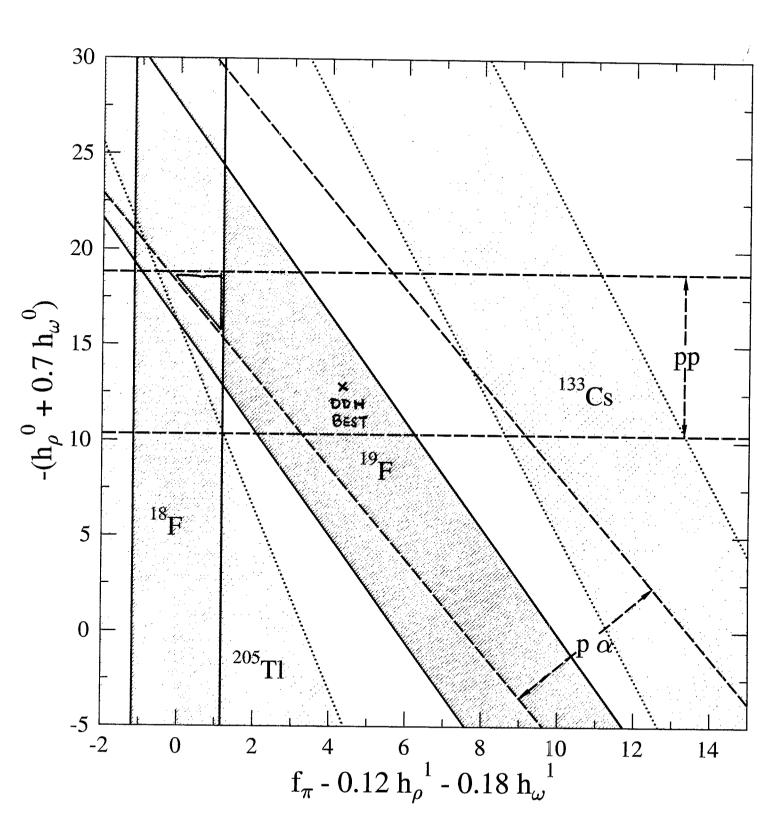
	$oldsymbol{A}$	В	C	D	E
¹³³ Cs	0.09	(0.08, 0.07)	0.05	0.04	0.05
²⁰⁵ Tl	-0.36	(-0.32, -0.30)	-0.28	-0.18	-0.08

■ Detailed breakdown of various contributions to AM

Nucleus	Source	f_{π}	$h_{ ho}^0$	$h_{ ho}^{1}$	$h_{ ho}^2$	h_{ω}^{0}	h^1_ω
133Cs	1 – body	0.59	0.87	0.90	0.36	0.28	0.29
	2 – body	8.58	0.02	0.11	0.06	-0.57	-0.57
	mixing	51.57	-16.67	-4.88	-0.06	-9.79	-4.59
	total	60.74	-15.78	-3.87	0.36	-10.09	-4.87
²⁰⁵ Tl	1 – body	-0.63	-0.86	-0.96	-0.35	-0.29	-0.29
· · · · · · · · · · · · · · · · · · ·	2 – body	-3.54	-0.01	-0.06	-0.03	0.28	0.28
	mixing	-13.68	4.63	1.34	0.08	2.77	1.27
•	total	-18.03	3.76	0.33	-0.30	2.76	1.26

■ Constraints on PNC meson-nucleon couplings

Observable	$Exp.(\times 10^{-7})$	$f_{\pi} - 0.12 h_{\rho}^{1}$	h_{ρ}^{0} + 0.7 h_{ω}^{0}	$\Delta h_{ ho}^{1}$	$h^2_{ ho}$	Δh_{ω}^{0}	Δh_{ω}^{1}
		$-0.18\mathrm{h}_\omega^1$		-			
A ^{pp} _Z (13.6 MeV)	-0.93 ± 0.21		0.043	0.043	0.017	0.009	0.039
A <mark>pp</mark> (45 MeV)	-1.57 ± 0.23		0.079	0.079	0.032	0.018	0.073
AZP (221 MeV)	0.84 ± 0.29		-0.030	-0.030	-0.012	0.021	
${ m A}_{ m Z}^{ m plpha}$ (46 MeV)	-3.34 ± 0.93	-0.340	0.140	0.006		-0.039	-0.002
P_{γ} (18F, 1081 keV)	1200 ± 3860	4385		34			-44
A _γ (¹⁹ F, 110 keV)	-740 ± 190	-94.2	34.1	-1.1		-4.5	-0.1
$\langle a (Cs) \rangle M_N^2 / e$	800 ± 140	60.7	-15.8	3.4	0.4	1.0	6.1
$\langle \parallel a (Tl) \parallel \rangle M_N^2 / e$	370 ± 390	-18.0	3.8	-1.8	-0.3	0.1	-2.0



Summary

- → Nuclear anapole moments are a sign of nuclear parity nonconservation and capable of putting good constraints on the PNC nuclear interaction. This means they provide a way to study the hadronic weak interaction and dynamics of strong interaction.
- → The theory-vs.-experiment discrepancy is one puzzle needed to be sorted out. More experiments alike are valuable, and better atomic and nuclear calculations are the challenge for theorists.